

Week 8 Worksheet Solutions

(Nondegenerate) Perturbation Theory

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October 18, 2024

Exercise 1. Let $H = H^0 + H^1$ be a perturbed hamiltonian. Suppose we know

$$H^0 |n^0\rangle = E_n^0 |n^0\rangle,$$

where the $|n^0\rangle$ are the unperturbed, orthonormal, nondegenerate eigenstates.

- Expand the exact solutions for H , $|n\rangle$ and E_n , in perturbation expansions.
- Write the Schrödinger equation for H in terms of the above expansions.
- Study the first order part of the equation from (b), and derive the first order corrections to the energies. You should get

$$E_n^1 = \langle n^0 | H^1 | n^0 \rangle.$$

- Along the way to solving (c), you should have come up with the equation

$$H^0 |n^1\rangle + H^1 |n^0\rangle = E_n^1 |n^0\rangle + E_n^0 |n^1\rangle.$$

Using this equation, find the expansion of $|n^1\rangle$ in the eigenbasis of H^0 .

Hints: In order to find the component of $|n^1\rangle$ which is parallel to $|n^0\rangle$, enforce normalization of $|n\rangle$ to first order, i.e. $|n^0\rangle + |n^1\rangle$ should have norm 1. It will be helpful to write $|n^1\rangle = |n_{\parallel}\rangle + |n_{\perp}\rangle$, where $\langle n^0 | n_{\perp} \rangle = 0$; also, use the fact that—to first order— $e^{ia} = 1 + ia$.

- Derive the second order corrections to the energies, E_n^2 .

- We write

$$|n\rangle = \sum_{k=0}^{\infty} \lambda^k |n^k\rangle$$
$$E_n = \sum_{k=0}^{\infty} \lambda^k E_n^k.$$

b) Recall that $H = H^0 + \lambda H^1$. Then we have

$$H |n\rangle = E_n |n\rangle$$

$$(H^0 + \lambda H^1) \sum \lambda^k |n^k\rangle = \sum \lambda^{k+j} E_n^k |n^j\rangle.$$

c) The order λ part of the equation from (b) is

$$H^0 |n^1\rangle + H^1 |n^0\rangle = E_n^1 |n^0\rangle + E_n^0 |n^1\rangle.$$

Bracket this equation with $\langle n^0|$ to obtain

$$E_n^0 \langle n^0 | n^1 \rangle + \langle n^0 | H^1 | n^0 \rangle = E_n^1 + E_n^0 \langle n^0 | n^1 \rangle,$$

from which we immediately obtain the desired result.

d) The idea is that we want to write $|n^1\rangle$ in the eigenbasis of H^0 . So we want to write

$$|n^1\rangle = \sum c_m |m^0\rangle,$$

and we know that the coefficients c_m are given by

$$c_m = \langle m^0 | n^1 \rangle.$$

Thus, we want to bracket the equation from (d) with $\langle m^0|$ with m not necessarily equal to n , and try to write down an equation for c_m . Doing this bracket, we have

$$E_m^0 \langle m^0 | n^1 \rangle + \langle m^0 | H^1 | n^0 \rangle = E_n^1 \delta_{nm} + E_n^0 \langle m^0 | n^1 \rangle.$$

Now, notice that if $m = n$, the equation we obtained doesn't tell us anything. But, if $m \neq n$, we immediately have

$$c_m = \frac{\langle m^0 | H^1 | n^0 \rangle}{E_n^0 - E_m^0}.$$

For the case $m = n$, we follow the hint. Normalization of $|n\rangle$ to first order gives

$$1 = \langle n^0 | n^0 \rangle + \langle n^0 | n_{\parallel}^1 \rangle + \langle n_{\parallel}^1 | n^0 \rangle,$$

where we note that terms like $\langle n^1 | n^1 \rangle$ are second order and can be ignored. Since $|n^0\rangle$ is also assumed normalized, we have

$$0 = 2\text{Re} [\langle n^0 | n_{\parallel}^1 \rangle],$$

so that we can set

$$|n_{\parallel}^1\rangle = i a |n^0\rangle,$$

with $a \in \mathbb{R}$. Thus,

$$\begin{aligned} |n\rangle &= (1 + ia) |n^0\rangle + |n^1_\perp\rangle \\ &= e^{ia} |n^0\rangle + |n^1_\perp\rangle. \end{aligned}$$

Since we can absorb the phase into the definition of $|n^0\rangle$, we find

$$\begin{aligned} |n^1\rangle &= |n^1_\perp\rangle \\ &= \sum_{m \neq n} \frac{\langle m^0 | H^1 | n^0 \rangle}{E_n^0 - E_m^0} |m^0\rangle, \end{aligned}$$

i.e. we can set $c_n = 0$.

e) We write down the second order equation obtained from the result of (b):

$$H^0 |n^2\rangle + H^1 |n^1\rangle = E_n^2 |n^0\rangle + E_n^1 |n^1\rangle + E_n^0 |n^2\rangle.$$

Then we bracket with $\langle n^0 |$ to find

$$\langle n^0 | H^1 | n^1 \rangle = E_n^2,$$

so

$$E_n^2 = \sum_{m \neq n} \frac{|\langle m | H^1 | n \rangle|^2}{E_n^0 - E_m^0}.$$

Exercise 2. Suppose you want to calculate the expectation value of some observable A in the n^{th} energy eigenstate of a system perturbed by H^1 ,

$$\langle A \rangle = \langle n | A | n \rangle.$$

Suppose further that all eigenstates are nondegenerate.

- Replace $|n\rangle$ by its perturbation expansion, and write down the formula for the first order correction to $\langle A \rangle$, $\langle A \rangle^1$.
- Use the first order corrections to the states,

$$|n^1\rangle = \sum_{m \neq n} \frac{\langle m^0 | H^1 | n^0 \rangle}{E_n^0 - E_m^0} |m^0\rangle,$$

to rewrite $\langle A \rangle^1$ in terms of the unperturbed eigenstates.

- If $A = H^1$, what does the result of (b) tell you? Explain why this is consistent with the result of Exercise 1(e),

$$E_n^2 = \sum_{m \neq n} \frac{|\langle m | H^1 | n \rangle|^2}{E_n^0 - E_m^0}.$$

a) We write

$$|n\rangle = |n^0\rangle + \lambda |n^1\rangle + \lambda^2 |n^2\rangle + \dots .$$

Thus,

$$\langle A \rangle = \langle n^0 | A | n^0 \rangle + 2\text{Re} \langle n^0 | A | n^1 \rangle + \dots ,$$

$$\text{so } \langle A \rangle^1 = 2\text{Re} \langle n^0 | A | n^1 \rangle .$$

b) Plugging in the expression for $|n^1\rangle$ given above, we get

$$\langle A \rangle^1 = 2\text{Re} \sum_{m \neq n} \frac{\langle m^0 | H^1 | n^0 \rangle}{E_n^0 - E_m^0} \langle n^0 | A | m^0 \rangle .$$

c) If $A = H^1$, then we get that the first order correction to the expectation value of H^1 is given by

$$2 \sum_{m \neq n} \frac{|H_{mn}^1|^2}{E_n^0 - E_m^0} ,$$

where $H_{mn}^1 = \langle m | H^1 | n \rangle$. On the other hand, the equation from 1(e) is

$$E_n^2 = \sum_{m \neq n} \frac{|H_{mn}^1|^2}{E_n^0 - E_m^0} .$$

So we have found

$$\langle H^1 \rangle^1 = 2E_n^2 .$$

On the other hand,

$$\begin{aligned} \langle H \rangle &= \langle H \rangle^0 + \langle H \rangle^1 + \langle H \rangle^2 \dots \\ &= \langle H^0 \rangle^0 + \langle H^1 \rangle^0 + \langle H^0 \rangle^1 + \langle H^1 \rangle^1 + \langle H^0 \rangle^2 + \dots , \end{aligned}$$

where the ellipsis denotes third and higher order terms. Thus, if we truncate to second order, we must have

$$E_n^2 = \langle H^1 \rangle^1 + \langle H^0 \rangle^2 ,$$

where $\langle H^0 \rangle^2$ denotes the expectation value of H^0 with respect to the state $|n^2\rangle$, and we expect this expectation value to be exactly $-E_n^2$.